

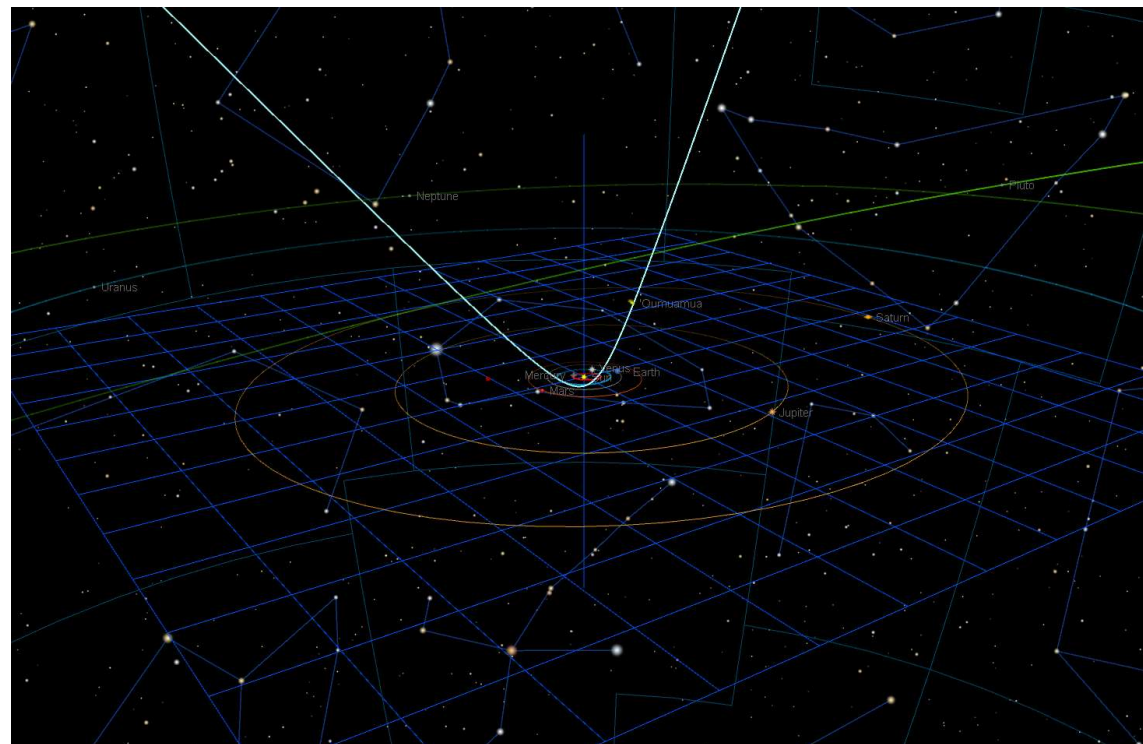
# WHAT ARE UNIVERSAL VARIABLES? *How Do They Describe the Path of the First-Known Interstellar Asteroid, 'Oumuamua?*

by Roger L. Mansfield, *MAA 50-Year Member*  
Astronomical Data Service, Colorado Springs, Colorado USA  
<http://astroger.com>

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*The figure above depicts the position of 'Oumuamua on its hyperbolic path, three months before its perihelion on 2017 September 9*

# Universal Variables Theories

- The eccentricity  $e$  of a two-body orbit or path indicates whether the motion is along an ellipse ( $e < 1$ ), a parabola ( $e = 1$ ), or a hyperbola ( $e > 1$ )
- In a “universal variables” theory, there is one set of equations that describes all of the possible conic paths of two-body motion
- Richard H. Battin [3], William H. Goodyear [4], and Samuel Herrick [5] all published universal variables theories of two-body motion in the 1960s
- But Karl J. Stumpff [1] had already published such a theory in 1947, using a new class of transcendental functions that have come to be called the Stumpff functions [6] or “c-functions”
- During this presentation, we will examine the vector equations for the propagation of a two-body path along an ellipse, a parabola, or a hyperbola, using a universal, two-body theory that traces back to Stumpff’s seminal article in the *Astronomische Nachrichten* in 1947

# Outline and Objective of this Presentation

- We view a depiction of the hyperbolic trajectory of the first-known interstellar asteroid, `Oumuamua, as an application of the vector equations of motion that we had just discussed
- We will then examine the first six c-functions in some detail
  - Series definitions and recurrence formulae
  - Derivatives and integrals
  - Quadruple-argument formulae
  - Calculation by series and recursion
- My objective for this presentation is that you will:
  - understand and appreciate the importance of Stumpff's c-functions in applied mathematics
  - understand how to calculate the first six c-functions
- This presentation is available at my website at <http://astroger.com>, along with a Borland C++ Builder 5 program that propagates the trajectory of `Oumuamua (and thus includes code that calculates the first six c-functions)

## Excerpt from Stumpff's "Neue Formeln" Article [1]

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K. STUMPF: Neue Formeln und Hilfstafeln zur Ephemeridenrechnung

und ist stets reell, da ja  $E$  und  $M$  stets gleichzeitig reell oder imaginär sind. Die Funktionen  $c_2$  und  $c_3$  entstammen der Folge

$$\left. \begin{aligned}
 c_0 &= \cos \lambda \\
 c_1 &= \frac{1}{\lambda} \int_0^\lambda \cos \lambda \, d\lambda = \frac{\sin \lambda}{\lambda} \\
 c_2 &= \frac{1}{\lambda^2} \int_0^\lambda \int_0^\lambda \cos \lambda \, (d\lambda)^2 = \frac{1 - \cos \lambda}{\lambda^2} \\
 \dots \\
 c_n &= \frac{1}{\lambda^n} \int_0^\lambda \dots \int_0^\lambda \cos \lambda \, (d\lambda)^n = \frac{1}{n!} \left[ 1 - \frac{\lambda^2}{(n+1)(n+2)} + \right. \\
 &\quad \left. + \frac{\lambda^4}{(n+1) \dots (n+4)} \dots \right]
 \end{aligned} \right\} \quad (10)$$

und sind als Funktionen von  $\lambda^2$  stets reell, auch wenn man für  $\lambda^2$  negative Werte (Hyperbeln) oder null (Parabeln) zuläßt.

## Second-Order ODE for Two-Body (Unperturbed) Orbital Motion

$$\frac{d^2\mathbf{r}}{dt^2} + K^2 \frac{\mathbf{r}}{r^3} = 0 \quad (\text{Here 3-vectors are boldface type})$$

For perturbed orbital motion,  $\frac{d^2\mathbf{r}}{dt^2} + K^2 \frac{\mathbf{r}}{r^3} = \mathbf{P}(\mathbf{r}, \frac{d\mathbf{r}}{dt})$ , where  $\mathbf{P}$  is a 3-vector of perturbative accelerations. In both equations now,

$\mathbf{r}$  (boldface) is position vector of secondary, with components  $x$ ,  $y$ , and  $z$

$\frac{d\mathbf{r}}{dt}$  is velocity vector of secondary, with components  $\frac{dx}{dt}$ ,  $\frac{dy}{dt}$ , and  $\frac{dz}{dt}$

$r$  (not boldface) is the magnitude of vector  $\mathbf{r}$ . That is,  $r = \sqrt{x^2 + y^2 + z^2}$

$K$  is gravity constant, i.e.,  $K = k \sqrt{1 + m/M}$ , where  $k$  is Gaussian constant for primary,  $m$  is mass of secondary and  $M$  is mass of primary

Note: Next slide uses overscript dot ( $\dot{\phantom{x}}$ ) instead of  $d/dt$  to denote derivative!

## Two-Body Orbit Propagation using Stumpff's c-Functions

Given  $\mathbf{r}_0$  and  $\dot{\mathbf{r}}_0$  at time  $t_0$ , find  $\mathbf{r}$  and  $\dot{\mathbf{r}}$  at  $t$  as follows.

Compute  $\alpha = 2/r_0 - (\dot{\mathbf{r}}_0 \cdot \dot{\mathbf{r}}_0)$ ;  $\sigma_0 = \mathbf{r}_0 \cdot \dot{\mathbf{r}}_0$ ;  $\tau = K(t-t_0)$ .

Then solve Stumpff's equation [1],

$$\tau = r_0 s c_1(\alpha s^2) + \sigma_0 s^2 c_2(\alpha s^2) + s^3 c_3(\alpha s^2),$$

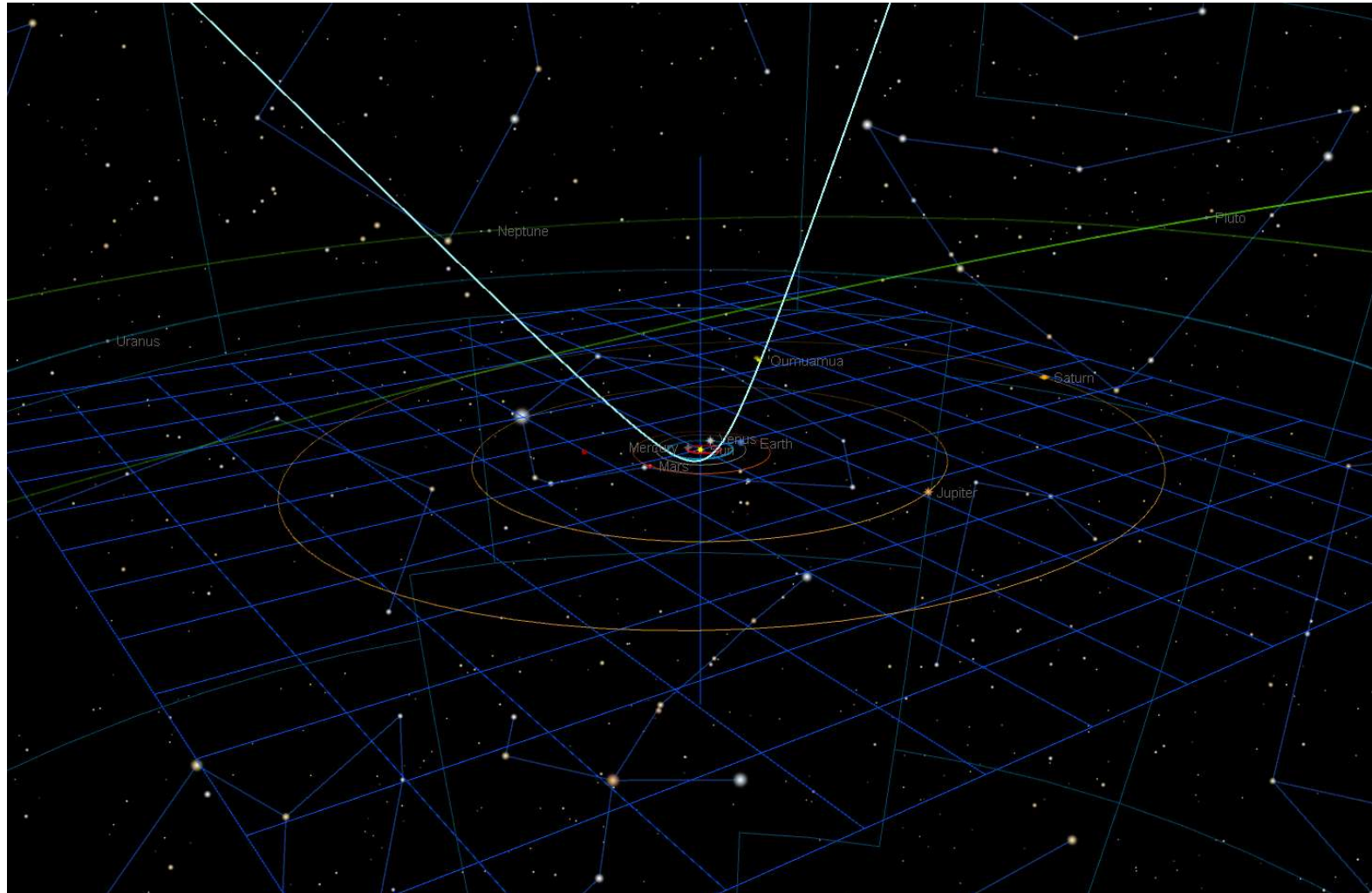
for  $s$  and  $r$ , using the fact that  $d\tau/ds = r$ . Finally,

$$\mathbf{f} = 1 - s^2 c_2(\alpha s^2)/r_0 \quad \mathbf{g} = \tau - s^3 c_3(\alpha s^2)$$

$$\dot{\mathbf{f}} = -s c_1(\alpha s^2)/(r r_0) \quad \dot{\mathbf{g}} = 1 - s^2 c_2(\alpha s^2)/r$$

$$\mathbf{r} = \mathbf{f} \mathbf{r}_0 + \mathbf{g} \dot{\mathbf{r}}_0 \quad \dot{\mathbf{r}} = \dot{\mathbf{f}} \mathbf{r}_0 + \dot{\mathbf{g}} \dot{\mathbf{r}}_0$$

# The Hyperbolic Trajectory of 'Oumuamua



*The figure above depicts the position of 'Oumuamua on its hyperbolic path, three months before its perihelion on 2017 September 9. It was created using Software Bisque's TheSky*



## Newton-Raphson Solution of Stumpff's Equation

First guess:  $s^{(k)} = \tau/r_0$ , for  $k = 0$

Iterate on:  $s^{(k+1)} = s^{(k)} - f(s^{(k)})/[df(s^{(k)})/ds]$

where  $f(s) = r_0 s c_1(\alpha s^2) + \sigma_0 s^2 c_2(\alpha s^2) + s^3 c_3(\alpha s^2) - \tau$

and  $df(s)/ds = r_0 c_0(\alpha s^2) + \sigma_0 s c_1(\alpha s^2) + s^2 c_2(\alpha s^2)$

$= r$

For Laguerre-Conway second-order solution [2], need second derivative,  $f''(s)$ :

$$f''(s) = dr/ds = -r_0 \alpha s c_1(\alpha s^2) + \sigma_0 c_0(\alpha s^2) + s c_1(\alpha s^2)$$



## Series Definitions of c-Functions

$$c_n(\lambda^2) = \sum_{k=0}^{\infty} (-1)^k \lambda^{2k} / (2k+n)! \quad (n = 0, 1, 2, \dots)$$

$$c_0(\lambda^2) = \frac{1}{0!} - \frac{\lambda^2}{2!} + \frac{\lambda^4}{4!} - \frac{\lambda^6}{6!} + \dots$$

$$c_1(\lambda^2) = \frac{1}{1!} - \frac{\lambda^2}{3!} + \frac{\lambda^4}{5!} - \frac{\lambda^6}{7!} + \dots$$

$$c_2(\lambda^2) = \frac{1}{2!} - \frac{\lambda^2}{4!} + \frac{\lambda^4}{6!} - \frac{\lambda^6}{8!} + \dots$$

$$c_3(\lambda^2) = \frac{1}{3!} - \frac{\lambda^2}{5!} + \frac{\lambda^4}{7!} - \frac{\lambda^6}{9!} + \dots$$

$$c_4(\lambda^2) = \frac{1}{4!} - \frac{\lambda^2}{6!} + \frac{\lambda^4}{8!} - \frac{\lambda^6}{10!} + \dots$$

$$c_5(\lambda^2) = \frac{1}{5!} - \frac{\lambda^2}{7!} + \frac{\lambda^4}{9!} - \frac{\lambda^6}{11!} + \dots$$

## Recurrence Formulae for c-Functions

$$c_n(\lambda^2) = \frac{1}{n!} - \lambda^2 c_{n+2}(\lambda^2)$$

for  $n = 0, 1, 2, \dots$

Follows by inspection, or formally, by "index manipulation" on the defining series.

Therefore, given  $c_5$  and  $c_4$  computed by series, it is possible to "proceed downward in  $n$ " to calculate  $c_3$ ,  $c_2$ ,  $c_1$ , and  $c_0$ .

Can also "proceed upward in  $n$ " from  $c_0$  and  $c_1$ . This permits greater efficiency of calculation in some cases.

## Integral Definitions of c-Functions

$$c_n(\lambda^2) = \frac{1}{\lambda^n} \int_0^\lambda \lambda^{n-1} c_{n-1}(\lambda^2) d\lambda \quad (n = 0, 1, 2, \dots)$$

For  $\lambda^2 \geq 0$

$$c_0(\lambda^2) = \cos \lambda$$

$$c_1(\lambda^2) = (\sin \lambda)/\lambda$$

$$c_2(\lambda^2) = (1 - \cos \lambda)/\lambda^2$$

$$c_3(\lambda^2) = (\lambda - \sin \lambda)/\lambda^3$$

For  $-\lambda^2 \leq 0$

$$c_0(-\lambda^2) = \cosh \lambda$$

$$c_1(-\lambda^2) = (\sinh \lambda)/\lambda$$

$$c_2(-\lambda^2) = (\cosh \lambda - 1)/\lambda^2$$

$$c_3(-\lambda^2) = (\sinh \lambda - \lambda)/\lambda^3$$

and for  $\lambda^2 = 0$ ,  $c_n(\lambda^2) = 1/n!$

## Differentiation of c-Functions

$$(1) \text{ For } n > 0, \quad \frac{d\lambda^n c_n(\lambda^2)}{d\lambda} = \lambda^{n-1} c_{n-1}(\lambda^2).$$

$$(2) \text{ For } n = 0, \quad \frac{dc_0(\lambda^2)}{d\lambda} = -\lambda c_1(\lambda^2) \quad (*).$$

Equation (2) follows from Eq. (1) and the recurrence formulae.

$$* \text{But note that } \frac{dc_0(\alpha s^2)}{ds} = -\alpha s c_1(\alpha s^2).$$

## Quadruple-Argument Formulae for c-Functions

The following formulae are needed in order to construct an efficient algorithm for calculation of  $c_0$ ,  $c_1$ ,  $c_2$ ,  $c_3$ ,  $c_4$ , and  $c_5$ .

$$c_0(4\lambda^2) = 2c_0(\lambda^2)c_0(\lambda^2) - 1$$

$$c_1(4\lambda^2) = c_0(\lambda^2)c_1(\lambda^2)$$

$$c_2(4\lambda^2) = [c_1(\lambda^2)c_1(\lambda^2)]/2$$

$$c_3(4\lambda^2) = [c_1(\lambda^2)c_2(\lambda^2) + c_3(\lambda^2)]/4$$

$$c_4(4\lambda^2) = [c_2(\lambda^2)c_2(\lambda^2) + 2c_4(\lambda^2)]/8$$

$$c_5(4\lambda^2) = [c_2(\lambda^2)c_3(\lambda^2) + c_4(\lambda^2) + c_5(\lambda^2)]/16$$

## Calculation of c-Functions

Need  $c_0, c_1, c_2, c_3, c_4,$  and  $c_5$  to compute state transition matrix. Need an efficient procedure for calculating all six of these c-functions, valid for any real argument:

- Reduce argument of c-functions  $N$  times ( $N \geq 0$ ) by quartering, until reduced argument is acceptably small\*.
- Calculate  $c_5$  and  $c_4$  of reduced argument using series, carrying as many terms as needed for desired precision\*.
- Calculate  $c_3, c_2, c_1,$  and  $c_0$  of reduced argument via recurrence formulae.

\*For this application,  $|\lambda^2| < 0.1$  was "smallness criterion", and seven terms were used in series for  $c_4$  and  $c_5$ .



## Calculation of c-Functions, Continued

- Calculate c-functions of original argument by iteration of the quadruple-argument formulae:

Let  $c_0^{(N)}$ ,  $c_1^{(N)}$ ,  $c_2^{(N)}$ ,  $c_3^{(N)}$ ,  $c_4^{(N)}$ , and  $c_5^{(N)}$  denote c-functions calculated for reduced argument. Iterate quadruple-argument formulae

$$c_5^{(k-1)} = [c_2c_3 + c_4 + c_5]^{(k)}/16$$

$$c_4^{(k-1)} = [c_2c_2 + c_4 + c_4]^{(k)}/8$$

$$c_3^{(k-1)} = [c_1c_2 + c_3]^{(k)}/4 \quad c_2^{(k-1)} = [c_1c_1]^{(k)}/2$$

$$c_1^{(k-1)} = [c_0c_1]^{(k)} \quad c_0^{(k-1)} = 2[c_0c_0]^{(k)} - 1$$

for  $k = N, N-1, \dots, 0$ . Note that if  $N = 0$  to begin with, then iteration is not necessary.



## Calculation of c-Functions, Summary Algorithm

- a. Set  $N = 0$  and set  $y = x$ , where  $x$  is input argument.
- b. If  $|y| < 0.1$  then go to Step d, else go to Step c.
- c. Increment  $N$  by 1 & divide  $y$  by 4, then go to Step b.
- d. Calculate

$$c_5 = (1-y(1-y(1-y(1-y(1-y(1-y/272)/210)/156) \\ /110)/72)/42)/120;$$

$$c_4 = (1-y(1-y(1-y(1-y(1-y(1-y/240)/182)/132) \\ / 90)/56)/30)/ 24;$$

$$c_3 = 1/6 - yc_5;$$

$$c_2 = 1/2 - yc_4;$$

$$c_1 = 1 - yc_3;$$

$$c_0 = 1 - yc_2.$$

- e. If  $N = 0$  then exit, else go to Step f.
- f. Decrement  $N$  by 1, calculate "new" c-functions of quadruple argument from "old" c-functions of previous argument, then go to Step e.

## Summary of Presentation

- **“Universal variables” is an algorithm that propagates the position and velocity of a secondary moving around its primary, and**
  - **the algorithm does not branch to three different sets of equations based upon the value of the eccentricity ( $e < 1$ ,  $e = 1$ , or  $e > 1$ )**
- **Karl J. Stumpff published such a universal variables algorithm in 1947 using his c-functions**
- **Stumpff's c-functions generalize the sine and cosine functions**
  - **we now know how to calculate the first six:  $c_0$ ,  $c_1$ ,  $c_2$ ,  $c_3$ ,  $c_4$ , and  $c_5$**
- **The solar system path of the recently discovered interstellar asteroid `Oumuamua is hyperbolic ( $e > 1$ ), so:**
  - **we can choose to propagate its path using an algorithm that only works for hyperbolic paths, or**
  - **we can choose to use the more general "universal variables" equations that work for all conic paths**

## REFERENCES

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[5] Herrick, Samuel, *Universal Variables*, *Astronomical Journal*, Vol. 70, No. 4 (May 1965), pp. 309-315.

[6] Stiefel, E. L. and G. Scheifele, *Linear and Regular Celestial Mechanics*, Springer-Verlag (New York, 1971), Section 11 (see p. 43).